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Baryogenesis of the universe in cSMCS model plus iso-doublet vector quark

Neda Darvishi

*Faculty of Physics, University of Warsaw,
Pasteura 5, 02-093 Warsaw, Poland*

E-mail: neda.darvishi@fuw.edu.pl

ABSTRACT: CP violation of the SM is insufficient to explain the baryon asymmetry in the universe and therefore an additional source of CP violation is needed. Here the extension of the SM by a neutral complex scalar singlet with a nonzero vacuum expectation value (cSMCS) plus a heavy vector quark pair is considered. This model offers the spontaneous CP violation and proper description in the baryogenesis, it leads strong enough first-order electro-weak phase transition to suppress the baryon-violating sphaleron process.

KEYWORDS: Beyond Standard Model, CP violation, Heavy Quark Physics, Higgs Physics

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1 Introduction

As far as we know, our universe is dominated by matter. The baryon asymmetry of the universe (BAU) could be generated during the electroweak phase transition, as disused in numerous studies, e.g. [1–13], trying to identify the source of this asymmetry. According to the Sakharov conditions [14], in order to generate the BAU it is required to have, first, the violation of the baryon number, second, the violation of C and CP symmetries and third, the existence of a non equilibrium processes, see also the reference.¹ It has become apparent that the electroweak Standard Model (SM) is unable to account for the observed magnitude of the BAU because the amount of CP violation within the SM is not sufficient [20–22]. In the present work it has been assumed that the generation of BAU is provided by the model with a neutral complex scalar singlet χ , which accompanies the SM-like Higgs doublet Φ , and an iso-doublet vector quarks $V_L + V_R$. This kind of extension of the SM was discussed in the literature with various motivations, can be found in eg. [23–26, 26–32].

Here the potential with a softly broken global U(1) symmetry is considered, leading the model called cSMCS, also see [33–35]. The additional source of CP violation is provided by a neutral complex scalar singlet χ with non-zero expectation value. The issue of the CP violation due to a complex singlet with a complex expectation value has been previously discussed, see [34, 35]. In the presence of an iso-doublet vector quark and a complex singlet, the Yukawa Lagrangian acquires additional terms. While diagonalizing the quark

¹For reviews see refs. [15–19].

mass matrix the whole Lagrangian will be modified with new terms which are functions of the time-dependent phase (CP violating phase [36]). The appearance of these terms leads to the generation of a baryon asymmetry.

The content of this paper is as follows. In section 2 a general presentation of the SMCS model and its constrained version (cSMCS) investigated in the paper is given. In particular, the subsection 2.2 describes the conditions for the spontaneous CP violation in the model. In section 3, the necessary conditions for strong enough first order electroweak phase transition in the model will be verified. The generation of a baryon asymmetry through the mixing of the SM quark and vector quarks is discussed in the section 4. The section 5 contains our conclusion while detailed formula are presented in the appendix.

2 The cSMCS: the SM plus a complex singlet scalar

The full Lagrangian of this model is given by

$$\mathcal{L} = \mathcal{L}_{gf}^{\text{SM}} + \mathcal{L}_{\text{scalar}} + \mathcal{L}_Y(\psi_f, \Phi) + \mathcal{L}_Y(V_q, \chi), \quad (2.1)$$

where $\mathcal{L}_{gf}^{\text{SM}}$ describes the pure gauge boson terms as well the SM boson- SM fermion interaction, $\mathcal{L}_{\text{scalar}}$ describes the scalar sector of the model with one SU(2) doublet Φ and a neutral complex scalar (spinless) singlet χ . $\mathcal{L}_Y(\psi_f, \Phi)$ and $\mathcal{L}_Y(V_q, \chi)$ represent, respectively, the Yukawa interaction of Φ with SM fermions and the Yukawa interaction of singlet scalar with vector quarks. In the cSMCS model the neutral singlet χ does not couple to the SM fermions and therefore the singlet- SM fermion interaction is present only through the mixing of the singlet χ with the doublet Φ (the same holds for the singlet interaction with the gauge bosons), while by adding vector quark the neutral singlet χ can couple to the SM fermions through vector quark. The SM-like Higgs boson in the model predominantly consists of a neutral CP-even component of the Φ doublet and its mass is ~ 125 GeV. There are two other higgs particles-like, see discussion in [33–35].

We assume Φ and χ fields have non-zero vacuum expectation values (vev) v and $we^{i\xi}$, respectively ($v, w, \xi \in \mathbf{R}$). We shall use the following field decomposition around the vacuum state,

$$\Phi = \begin{pmatrix} \phi^+ \\ \frac{1}{\sqrt{2}}(v + \phi_1 + i\phi_4) \end{pmatrix}, \chi = \frac{1}{\sqrt{2}}(we^{i\xi} + \phi_2 + i\phi_3), \quad (2.2)$$

where

$$we^{i\xi} = w \cos \xi + iw \sin \xi = w_1 + iw_2. \quad (2.3)$$

Masses of the gauge bosons are given by the vev of the doublet, e.g $M_W^2 = g^2 v^2 / 4$ for the W boson.

2.1 Potential

The scalar potential of the model can be written as follows [33–35]

$$V = V_D + V_S + V_{DS}, \quad (2.4)$$

with the pure doublet and the pure singlet parts (respectively V_D and V_S) and the mixed term V_{DS} . The SM part of the potential represent by V_D , is equal to

$$V_D = -\frac{1}{2}m_{11}^2\Phi^\dagger\Phi + \frac{1}{2}\lambda_1\left(\Phi^\dagger\Phi\right)^2. \quad (2.5)$$

The potential for a complex singlet is equal to,

$$V_S = -\frac{1}{2}m_s^2\chi^*\chi - \frac{1}{2}m_4^2(\chi^{*2} + \chi^2) + \lambda_{s1}(\chi^*\chi)^2 + \lambda_{s2}(\chi^*\chi)(\chi^{*2} + \chi^2) + \lambda_{s3}(\chi^4 + \chi^{*4}) \\ + \kappa_1(\chi + \chi^*) + \kappa_2(\chi^3 + \chi^{*3}) + \kappa_3(\chi + \chi^*)(\chi^*\chi). \quad (2.6)$$

The doublet-singlet interaction term is,

$$V_{DS} = \Lambda_1(\Phi^\dagger\Phi)(\chi^*\chi) + \Lambda_2(\Phi^\dagger\Phi)(\chi^{*2} + \chi^2) + \kappa_4(\Phi^\dagger\Phi)(\chi + \chi^*). \quad (2.7)$$

There are three quadratic (m_i^2), six dimensionless quartic (λ_i, Λ_i) and four dimensionful parameters κ_i , $i = 1, 2, 3, 4$, describing linear (κ_1) and cubic terms (κ_2, κ_3) and κ_4 . The linear term κ_1 can be removed by a translation of the singlet field. Both V_S and V_{DS} are symmetry under the $\chi \rightarrow \chi^*$ transformation. To simplify the model, we apply a global U(1) symmetry [33–35].

$$U(1) : \quad \Phi \rightarrow \Phi, \chi \rightarrow e^{i\alpha}\chi. \quad (2.8)$$

However, a non-zero *vev* of χ would lead in such case to a spontaneous breaking of this symmetry and an appearance of massless Nambu-Goldstone scalar particles, what is not acceptable. Keeping some U(1) soft-breaking terms in the potential would solve this problem. In what follows, we shall consider a potential with a soft-breaking of U(1) symmetry, where the singlet cubic terms $\kappa_{2,3}$, κ_4 and the singlet quadratic term m_4^2 are kept. Therefore, we are left with the U(1)-symmetric terms ($m_{11}^2, m_s^2, \lambda_1, \lambda_{s1}, \Lambda_1$) and the U(1)-soft-breaking terms ($m_4^2, \kappa_{2,3,4}$). Simplifying slightly the notation by using: $\lambda_s = \lambda_{s1}, \Lambda = \Lambda_1$, we get the potential in the following form

$$V = -\frac{1}{2}m_{11}^2\Phi^\dagger\Phi + \frac{1}{2}\lambda_1\left(\Phi^\dagger\Phi\right)^2 + \Lambda(\Phi^\dagger\Phi)(\chi^*\chi) - \frac{1}{2}m_s^2\chi^*\chi + \lambda_s(\chi^*\chi)^2 + \kappa_4(\Phi^\dagger\Phi)(\chi + \chi^*) \\ - \frac{1}{2}m_4^2(\chi^{*2} + \chi^2) + \kappa_2(\chi^3 + \chi^{*3}) + \kappa_3(\chi + \chi^*)(\chi^*\chi). \quad (2.9)$$

Note that V is symmetry under the $\chi \rightarrow \chi^*$ transformation and we take all parameters real. Therefor V is explicitly CP conserving. We shall call the model with this choice of parameters, cSMCS [34, 35]. Note, that this potential (2.9) is similar to the potential with two real singlets, with an additional Z_2 symmetry for the one singlet field, considered in paper [31]. In that model, however, CP violation is not possible.

The extremum conditions lead to the following constraints,

$$-m_{11}^2 + v^2\lambda_1 + 2\sqrt{2}w_1\kappa_4 + \Lambda w^2 = 0, \quad (2.10)$$

$$w_1(-\mu_1^2 + v^2\Lambda + 2w^2\lambda_s) + \sqrt{2}[3(w_1^2 - w_2^2)\kappa_2 + (3w_1^2 + w_2^2)\kappa_3] + v^2\sqrt{2}\kappa_4 = 0, \quad (2.11)$$

$$w_2[-\mu_2^2 + v^2\Lambda + 2w^2\lambda_s + 2\sqrt{2}w_1(-3\kappa_2 + \kappa_3)] = 0, \quad (2.12)$$

Where the parameters μ_1^2 and μ_2^2 defined as

$$\mu_1^2 = m_s^2 + 2m_4^2, \quad \mu_2^2 = m_s^2 - 2m_4^2.$$

Various spontaneous symmetry breaking extrema are possible, among them with vanishing one or two of vacuum expectation parameters v, w_1, w_2 . Here we concentrate on the case with v, w_1 and w_2 different from zero, allowing for a vacuum which violate CP. Note, that for vanishing phase $\xi = 0$ ($w_2 = 0$) CP violation is not possible for our model [34].

In order to have a stable minimum, the parameters of the potential need to satisfy the positivity conditions. Positivity conditions read as follows:

$$\lambda_1, \lambda_s > 0, \quad \Lambda > -\sqrt{2\lambda_1\lambda_s}. \quad (2.13)$$

In addition the unitarity condition reaches to limit on quartic terms λ_1, λ_{s1} and Λ , below 4π .

2.2 The CP violating vacuum

Minding the equation (2.10), when neither v, w_1 nor w_2 vanish, an important relation can be obtained via subtracting equation (2.12) from the equation (2.11),

$$-8m_4^2 \cos^2 \xi + 6R_2 \cos \xi (1 + 2 \cos 2\xi) + 2R_3 \cos \xi + R_4 = 0, \quad (2.14)$$

where

$$R_2 = \sqrt{2}w\kappa_2, \quad R_3 = \sqrt{2}w\kappa_3, \quad R_4 = \frac{2\sqrt{2}v^2\kappa_4}{w} \cos \xi,$$

all of which are of $[mass]^2$ dimension. In addition we have

$$R_4 = \frac{v^2}{w^2}(m_{11}^2 - v^2\lambda_1 - w^2\Lambda).$$

For a particular case, i.e. $R_2 = 0$, the above equation transforms to,

$$-8m_4^2 \cos^2 \xi + 2R_3 \cos \xi + R_4 = 0. \quad (2.15)$$

In figure 1 the regions allowed by equations (2.10), (2.11) and (2.12) of parameters for a vacuum with $v, w_1, w_2 \neq 0$ is presented. Figure 1(a) shows the region of parameters R_3, R_4 and ξ as given by the eq. (2.15), for fixed m_4^2 is shown.

In figure 1(b) and (c) the allowed regions of parameters (R_3, R_4) , for $R_3 = 0$, and the allowed region of the parameters (R_2, R_4) , for $R_3 = 0$ are shown, respectively. These regions are in agreement with the eq. (2.14), for fixed m_4^2 .

The various aspect of CP violation with $\kappa_4 = 0$, are discussed in [34].

3 The electroweak phase transition

At very high temperatures far above the electroweak scale, the electroweak gauge symmetry is unbroken with no baryon number. The universe cools down and expands. Near the electroweak phase transition(EWPT) scale bubbles of broken electroweak symmetry appear and subsequently expand in the surrounding unbroken phase. Electroweak baryon

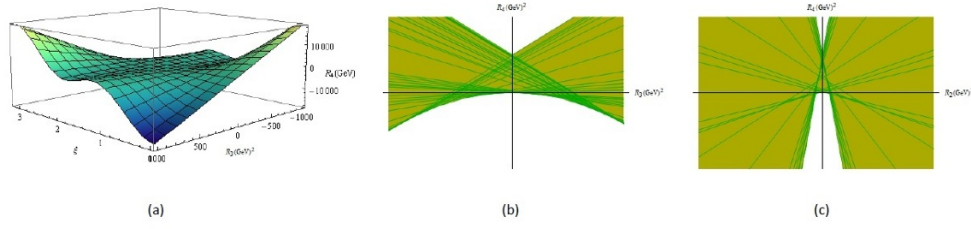


Figure 1. The spontaneously CP violation: the allowed (shaded) regions of the parameters R_2, R_3, R_4 for $-1 < \cos \xi < 1$ and $4m_4^2 = 500 \text{ GeV}^2$. The boarder lines correspond to the $\cos \xi = \pm 1$ limits. (a) Plot 3d for parameters R_3, R_4 and ξ based on eq. (2.15); (b) Regions for R_3, R_4 allowed by eq. (2.14); (c) Regions for R_2 and R_4 given by eq. (2.14).

asymmetry can be realized if this change of phase proceeds by a first-order phase transition. If the phase transition is strongly first order, the baryon violating processes are out of equilibrium in the bubble walls and a net baryon number can be generated during the phase transition. Phase transition is strong enough when [38–40],

$$\frac{v(T_c)}{T_c} \geq 1, \quad (3.1)$$

where T_c corresponds to the critical temperature.

To study the EWPT in the present model, we express the complex scalar χ in terms of its real and imaginary parts, $\chi = (\chi_1 + i\chi_2)/\sqrt{2}$. For the potential at zero temperature we have

$$V(T_0) = -\frac{1}{2}m_{11}^2\Phi^\dagger\Phi + \frac{1}{2}\lambda_1(\Phi^\dagger\Phi)^2 - \frac{\mu_1^2}{4}\chi_1^2 - \frac{\mu_2^2}{4}\chi_2^2 + \frac{1}{2}\Lambda(\Phi^\dagger\Phi)(\chi_1^2 + \chi_2^2) + \frac{1}{4}\lambda_s(\chi_1^2 + \chi_2^2)^2 + \frac{1}{\sqrt{2}}\kappa_2(\chi_1^3 - 3\chi_1\chi_2^2) + \frac{1}{\sqrt{2}}\kappa_3(\chi_1^3 + \chi_1\chi_2^2) + \sqrt{2}\kappa_4(\Phi^\dagger\Phi)\chi_1. \quad (3.2)$$

The one-loop thermal corrections to the effective potential at finite temperature T are (see ref. [40] for review),

$$\Delta V_{\text{thermal}} = \sum_i \frac{n_i T^4}{2\pi^2} I_{B,F} \left(\frac{m_i^2}{T^2} \right), \quad (3.3)$$

with

$$I_{B,F}(y) = \int_0^\infty dx x^2 \ln \left[1 \mp e^{-\sqrt{x^2+y}} \right]. \quad (3.4)$$

The minus sign correspond to the bosons and the plus to sign for the fermions. In eq. (3.3) m_i is the field-dependent mass and n_i is the number of degrees of freedom, (see appendix eq. (A.1)). One-loop potential for our model, calculated using the high temperature approximation (i.e. keeping only T^2 terms) is as follows [41],

$$V(T) = \frac{1}{4}\bar{m}_{11}^2\phi_1^2 + \frac{1}{8}\lambda_1\phi_1^4 + \frac{\bar{\mu}_1^2}{4}\phi_2^2 + \frac{\bar{\mu}_2^2}{4}\phi_3^2 + \frac{1}{4}\Lambda\phi_1^2(\phi_2^2 + \phi_3^2) + \frac{1}{4}\lambda_s(\phi_2^2 + \phi_3^2)^2 + \kappa_2\frac{1}{\sqrt{2}}(\phi_2^3 - 3\phi_2\phi_3^2) + \kappa_3\frac{1}{\sqrt{2}}(\phi_2^3 + \phi_2\phi_3^2) + \frac{1}{\sqrt{2}}\kappa_4\phi_1^2\phi_2 + \bar{\kappa}_{34}\frac{T^2}{3}\phi_2, \quad (3.5)$$

where

$$\begin{aligned}
 \bar{m}_{11}^2 &= -m_{11}^2 + \left(3\lambda_1 + \Lambda + \frac{2m_W^2 + m_Z^2 + 2m_t^2}{2v^2} \right) \frac{T^2}{3}, \\
 \frac{1}{2}\bar{\mu}_1^2 &= -\frac{1}{2}\mu_1^2 + (\Lambda + 2\lambda_s) \frac{T^2}{3}, \\
 \frac{1}{2}\bar{\mu}_2^2 &= -\frac{1}{2}\mu_2^2 + (\Lambda + 2\lambda_s) \frac{T^2}{3}, \\
 \bar{\kappa}_{34} &= \sqrt{2}(\kappa_3 + \kappa_4).
 \end{aligned} \tag{3.6}$$

The extremum conditions of the effective potential at temperature T , with respect to the fields ϕ_1, ϕ_2 and ϕ_3 , are

$$\bar{m}_{11}^2 + \lambda_1 v^2 + \Lambda w^2 + 2\sqrt{2}\kappa_4 w_1 = 0, \tag{3.7}$$

$$w_1(\bar{\mu}_1^2 + \Lambda v^2 + 2\lambda_s w^2) + \sqrt{2} [3\kappa_2(w_1^2 - w_2^2) + \kappa_3(3w_1^2 + w_2^2) + \kappa_4 v^2] + \frac{2}{3}\bar{\kappa}_{34}T^2 = 0, \tag{3.8}$$

$$w_2[\bar{\mu}_2^2 + \Lambda v^2 + 2\lambda_s w^2 + 2\sqrt{2}(-3\kappa_2 + \kappa_3)w_1] = 0. \tag{3.9}$$

At extreme temperatures, the solution of the equations (3.8), (3.8) and (3.9) is

$$v = 0, \quad w_1 \approx \frac{\bar{\kappa}_{34}}{2\lambda_s + \Lambda}, \quad w_2 \approx 0. \tag{3.10}$$

Then, the scalar component, ϕ_3 , decouples from the model and therefore, the potential is similar to the SM plus a real singlet [42].

Now, we will verify the parameter region where exists the strong first order phase transition, $v(T_c)/T_c \geq 1$, with the critical temperature T_c to be smaller than 250 GeV and $v(T_c)$ below its zero-temperature value $v_0 = 246$ GeV. We have performed a scan over the parameter space fulfilling unitarity and positivity conditions (see [34]), namely:

$$\begin{aligned}
 -0.25 < \Lambda < 0.25, \quad 0 < \lambda_s < 1, \quad -1 < \rho_{2,3,4} < 1, \quad 0 < \xi < \pi, \\
 -90000 \text{ GeV}^2 < \mu_1^2, \mu_2^2, \quad m_{11}^2 < 90000 \text{ GeV}^2,
 \end{aligned} \tag{3.11}$$

where we used dimensionless parameters $\rho_{2,3,4} = \kappa_{2,3,4}/w$. Considering the SM-like scenarios at the LHC, the mass of lightest higgs boson in this model is given by $M_{h_1}^2 \approx m_{11}^2 \approx \lambda_1 v^2$ ($M_{h_1} \approx 125$ GeV). Thus, we take λ_1 in the range [34]:

$$0.2 < \lambda_1 < 0.3. \tag{3.12}$$

The model contains two additional higgs scalars M_{h_2} and M_{h_3} , which we take to be [33, 34]

$$M_{h_3} \gtrsim M_{h_2} > 150 \text{ GeV}. \tag{3.13}$$

In [34] we have shown that these ranges of parameters are in agreement with LHC data and measurements of the oblique parameters.

The results of our scan are shown in the figure 2. In figure 2(a) the allowed region of $v(T_c)/T_c$ as a function of T_c is shown. Within the interval $100 < T_c < 200$ the ratio $v(T_c)/T_c$ can reach 2.5. The figure 2(b) shows that $v(T_c)/T_c \geq 1$ is possible for $|\rho_3| > 10^{-3}$.

We see that the strongly first-ordered EWPT is possible in our model. Since the out of equilibrium condition can be achieved for strong enough first order phase transition, in the bubble walls, we conclude that a successful BAU in our model is possible [40].

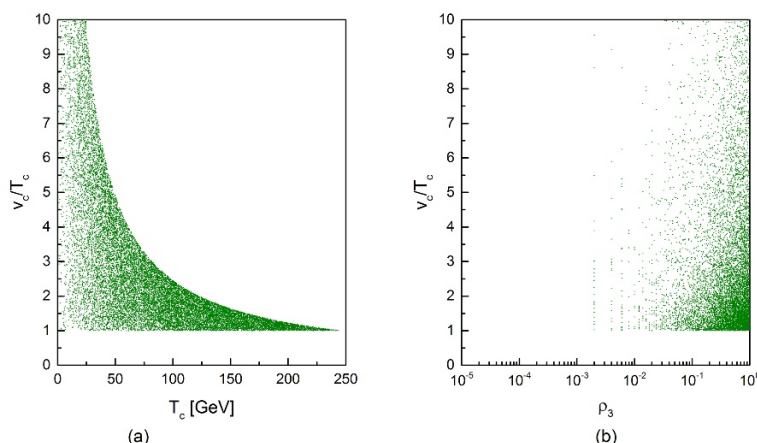


Figure 2. The allowed regions of critical temperature T_c , $v_c \equiv v(T_c)$ and $|\rho_3|$ for strongly first order phase transition. (a) $(T_c, v_c/T_c)$ and (b) $(|\rho_3|, v_c/T_c)$. The scatter points are selected to satisfy the criterion, $(v(T_c)/T_c) \geq 1$ (see text for details).

4 Baryogenesis

In this section we describe the baryon asymmetry resulting from a mixing of the SM quarks and heavy vector quarks [36, 42]. We consider a pair (Left and Right handed) of heavy iso-doublet vector quarks, $V_L + V_R$, with V_L and V_R having the same transformation properties under the gauge group of the SM, and where V_L transforms in the same way as a standard model quark doublet Q_L . Essentially the case of iso-doublet quarks have the same result to the case with iso-singlet vector quarks with transforming as either u_R or d_R with respect to the SM gauge group [36]. In order to have more general expressions and have both up and down types of vector quarks (respectively, \mathcal{U} and \mathcal{D}), is preferred to utilize the iso-doublet vector quarks. The transformation of the fermions under the additional global U(1) symmetry is as follows:

$$\text{U}(1) : \psi \rightarrow e^{i\alpha} \psi. \quad (4.1)$$

under this symmetry, left and right handed components rotate with the same phase so $\bar{\psi}\psi$ remains unchanged. The mass terms in the presence of the complex singlet are (see eq. (2.1)):

$$\mathcal{L}_Y(V_q, \chi) = \lambda_V \chi \bar{Q}_L V_R + M \bar{V}_L V_R + h.c. \quad (4.2)$$

In general the couplings of all three generations of SM quarks to vector quarks should be considered, nevertheless, since only one linear combination of these will actually mix with the vector quarks, we consider only one (the heaviest) quark Q_L . To generate baryon asymmetry, the phase of the singlet vev should be time-dependent, otherwise, such constant phase can be easily rotated away with the redefining the V_L and V_R [36]. Diagonalizing the quark mass matrix results in some non-diagonal kinetic terms. In addition, couple of time-dependent terms appear in the Lagrangian (see appendix eq. (C.1)), namely

$$\bar{Q}_L i\gamma^\mu \partial_\mu Q_L + \bar{V}_L i\gamma^\mu \partial_\mu V_L \rightarrow \bar{Q}'_L i\gamma^\mu \partial_\mu Q'_L + \bar{V}'_L i\gamma^\mu \partial_\mu V'_L + \Delta\mathcal{L}_k + \text{const}. \quad (4.3)$$

Since the CP violation disappears for a constant phase, when calculating the baryon asymmetry only the following kinetic term needs to be considered

$$\Delta\mathcal{L}_k = -\frac{\lambda_V^2 w^2}{M^2} \dot{\xi} (\bar{Q}'_L \gamma^0 Q'_L - \bar{V}'_L \gamma^0 V'_L). \quad (4.4)$$

Such term increases the baryon density of the universe with the transport of charge into the bubble wall. Conventionally, the amount of BAU is calculated via the following relation,

$$n_B = -N_f \int \frac{\Gamma_{\text{sph}}(T)}{2T} \mu_B dt, \quad (4.5)$$

where N_f is the number of flavors in the model. The sphaleron rate, Γ_{sph} , is defined as $\Gamma_{\text{sph}} = K(\alpha_W T)^4$ in the symmetric phase. K is the numerical factor reflecting the uncertainty in the estimate of the transition rate between vacua of different $B + L$ value, It has been estimated to be between 0.1 and 1 [43]. The chemical potential, μ_B , is associated with the baryonic charge in n_B . The chemical potential for third generation is as follows [36],

$$\mu_B = -\frac{5}{6} \frac{\lambda_V^2 w^2}{M^2} \dot{\xi}. \quad (4.6)$$

We assume that the mass parameter M is much larger than the temperature, i.e. $M \geq T$. Thus the sphaleron fluctuations cannot produce V quark pairs and we get the number density of baryons n_B at the temperature T as follows

$$n_B = \frac{5K\alpha_W^4}{2} \frac{\lambda_V^2 w^2}{M^2} \delta\xi T^3, \quad (4.7)$$

where $\delta\xi$ is the total change of the phase ξ . The BAU is determined via the ratio of the baryon number to the entropy [45]. The entropy density can be defined as

$$s = \frac{2\pi^2}{45} g^* T^3, \quad (4.8)$$

therefore,

$$\frac{n_B}{s} = \frac{225K\alpha_W^4}{4\pi^2 g^*} \frac{\lambda_V^2 w^2}{M^2} \delta\xi, \quad (4.9)$$

where the SU(2) gauge coupling $\alpha_W = 3.4 \times 10^{-2}$ and $g^* \sim 100$ is the effective number of degrees of freedom in the thermal equilibrium. The observations from WMAP gives the following value for n_B/s ratio [44, 46]

$$\frac{n_B}{s} = 8.7 \pm 0.3 \times 10^{-11}. \quad (4.10)$$

From eq. (4.9) and eq. (4.10) we get,

$$K \frac{\lambda_V^2 w^2}{M^2} \delta\xi = 1.14 \pm 0.3 \times 10^{-3}. \quad (4.11)$$

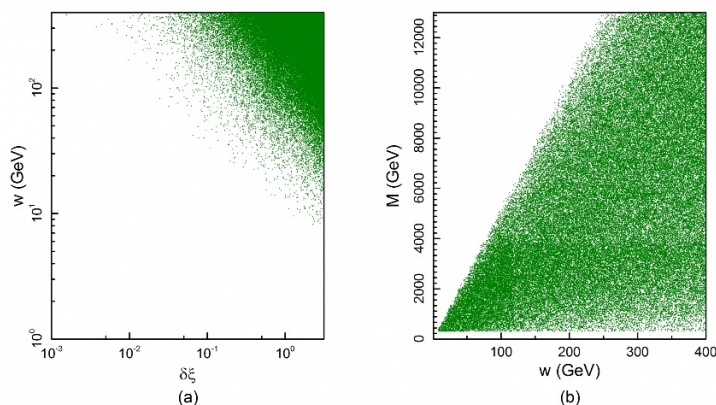


Figure 3. The allowed region for $\delta\xi$, w and M in the eq. (4.10) of the given ranges, i.e. eq. (4.12), for acceptable BAU within 2σ . (a) is the correlation of $\delta\xi$ and w . (b) is the correlation of w versus M .

The numerical analysis of the equations (4.9) and (4.10) have been performed via scanning the involving parameters in the following range,

$$\begin{aligned}
 0.3 \text{ TeV} &< M < 13 \text{ TeV}, \\
 0 &< \lambda_V < 1, \\
 2 \text{ GeV GeV} &< w < 400 \text{ GeV}, \\
 0 &< \delta\xi < \pi,
 \end{aligned}
 \tag{4.12}$$

with numerical factor $K = 1$. Figure 3 illustrates the parameter space allowing the generation of observed BAU for the n_B/s ratio within 2σ . The result were obtained from scanning in the ranges given by eq. (4.12) with the central value of eq. (4.10). Notice that the range of w is chosen in agreement with our previous analyses [34]. In figure 3(a) the distribution of the parameters $(\delta\xi, w)$ is shown. The parameter space for the (w, M) points is shown in the figure 3(b). Since M and w are independent parameters, their correlation is a direct consequence of the constraint (4.11). Based on these results, we conclude that our model provide successful BAU.

5 Conclusion

In the present work the possibility of the first-order EWPT for the cSMCS model is investigated, showing that such process is strong enough to generate BAU. In this analysis we have found that cubic terms, κ_2 , κ_3 and/or κ_4 of the potential, should have non-zero values. This is in agreement with our previous findings, regarding the parameter space of the possible region for CP violation in the cSMCS model, see the reference [34]. Afterwards, the parameter space of the model for the valid regions of BAU is scanned, concluding that the enlargement of the cSMCS model with a heavy iso-doublet vector quark pair could successfully predict an acceptable value for BAU.

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A Evaluation of the integral

The integral eq. (3.4) is evaluated as follow,

$$\frac{\partial}{\partial y} I_{B,F}(y) = \frac{1}{2} \int_0^\infty dx \frac{x^2}{(x^2 + y)^{1/2}} \frac{1}{\exp((x^2 + y)^{1/2}) - 1}, \quad (\text{A.1})$$

$$I_{B,F}(y)|_{y=0} = \int_0^\infty dx x^2 \ln(1 - e^{-x}) = -\frac{\pi^4}{45}, \quad (\text{A.2})$$

$$\frac{\partial}{\partial y} I_{B,F}(y)|_{y=0} = \frac{1}{2} \int_0^\infty dx \frac{x}{e^x - 1} = \frac{\pi^2}{12}. \quad (\text{A.3})$$

B The field-dependent mass m_i

The field-dependent mass m_i of gauge bosons, Goldstone boson, m_{ϕ_1} , m_{ϕ_2} and m_{ϕ_3} , which is used in eq. (3.3) are given by,

$$\begin{aligned} M_W^2 &= \frac{g^2 \phi_1^2}{4}, M_Z^2 = (g^2 + g'^2) \frac{\phi_1^2}{4}, \\ m_G^2 &= \lambda_1 \phi_1^2 + \Lambda(\phi_2^2 + \phi_3^2) + 2\sqrt{2}\kappa_4 \phi_2, \\ m_\phi^2 &= 3\lambda_1 \phi_1^2 + \Lambda(\phi_2^2 + \phi_3^2) + 2\sqrt{2}\kappa_4 \phi_2, \\ m_{\phi_2}^2 &= 3\lambda_s \phi_2^2 + \lambda_s \phi_3^2 + 3\sqrt{2}(\kappa_2 + \kappa_3)\phi_2 + \frac{1}{2}\Lambda\phi_1^2, \\ m_{\phi_3}^2 &= 3\lambda_s \phi_3^2 + \lambda_s \phi_2^2 + \sqrt{2}(-3\kappa_2 + \kappa_3)\phi_2 + \frac{1}{2}\Lambda\phi_1^2. \end{aligned} \quad (\text{B.1})$$

n_i is the number of degrees of freedom as,

$$n_W = 6, n_Z = 3, n_G = 3, n_{\phi, \phi_2, \phi_3} = 1, n_t = 12. \quad (\text{B.2})$$

C Rotation matrix

As discussed in section 4 the transformation of Q and V with rotation matrix read Q' and V' , shows as follows

$$\begin{aligned} \begin{bmatrix} Q'_L \\ V'_L \end{bmatrix} &= \begin{bmatrix} a & b \\ -b^* & a^* \end{bmatrix} \begin{bmatrix} Q_L \\ V_L \end{bmatrix} \\ a &= \left[1 + \left(\frac{\lambda_v w}{M} \right)^2 \right]^{-1/2} \\ b &= \left(\frac{\lambda_v w}{M} \right) \left[1 + \left(\frac{\lambda_v w}{M} \right)^2 \right]^{-1/2} e^{-i\xi} \end{aligned} \quad (\text{C.1})$$

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